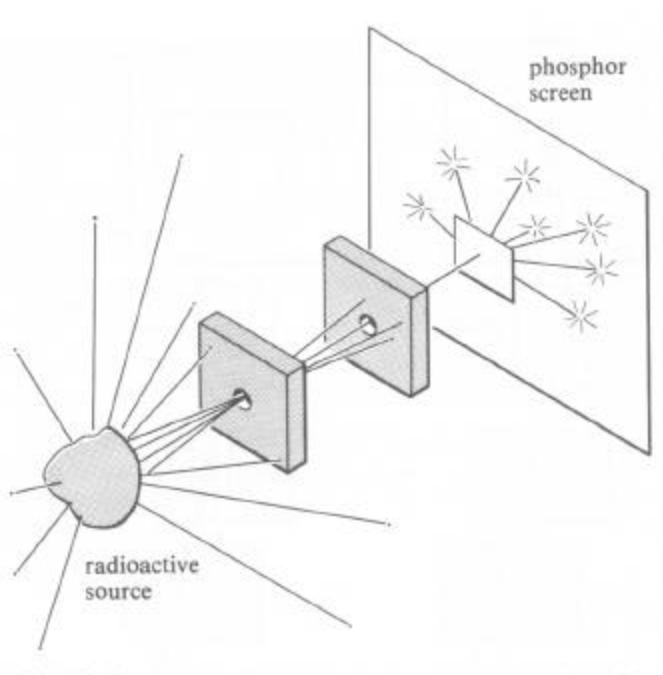


Atoms and Molecules

Rutherford (1910): scattering of charged particles by matter (gold foil).



Conclusion: Atoms consist of mostly empty space. Positively charged nucleus occupies very small volume in each atom.

Atomic Structure

Atoms: nucleus $\sim 10^{-14}$ m Z protons, charge $+e$

N neutrons, charge 0

electron cloud $\sim 10^{-10}$ m Z electrons, charge $-e$

Atomic number Z : 1 to 92 for naturally occurring elements

Atomic mass $A = Z + N$

$e = 1.602 \times 10^{-19}$ C

Z determines the element, A the isotope. Elements identified as ${}^A X$,

e.g., ${}^1\text{H}$, hydrogen; ${}^{12}\text{C}$, carbon; ${}^{16}\text{O}$, oxygen.

Atomic Structure - Classical Theory

Electric (Coulomb) force $F_e = \frac{kq_1q_2}{r^2} = -\frac{k(Ze)e}{r^2}$.

Energy ($KE + PE$) $E = \frac{mv^2}{2} - \frac{kZe^2}{r}$.

But, classical picture of orbiting electron means it accelerates.
Acceleration \Rightarrow radiation of EM waves \Rightarrow loss of energy \Rightarrow
collapse of orbit!

Wave-Particle Duality

de Broglie (1910's): proposes particles have wave-like properties.

$p = h/\lambda$, where p is the momentum of a particle.

Bohr (1913): a semi-classical model for the atom.

Only discrete electron orbit radii are allowed, and when in those orbits, the electron cannot radiate.

Allowed orbits for $mvr = n(h / 2\pi) \quad n = 1,2,3,\dots$

Note: along with deBroglie's hypothesis, implies orbit circumference must equal an integer # of wavelengths.

Combine with $\frac{mv^2}{r} = \frac{kZe^2}{r^2}$.

Bohr Model

Permitted levels $r = n^2 (h^2 / 4\mathbf{p}^2 me^2 kZ)$.

$$\text{Energy } E = \frac{mv^2}{2} - \frac{kZe^2}{r} \Rightarrow E(n) = -(2\mathbf{p}^2 me^4 k^2 Z^2) / n^2 h^2.$$

Orbits are bound until $n \rightarrow \infty$. Ground state $n = 1$.

Electrons can make a transition from one bound level to another. A photon is emitted or absorbed.

$$E(n_a) = E(n_b) + h\nu \quad (\text{emission}) \quad \text{where } n_a > n_b.$$

$$E(n_b) + h\nu = E(n_a) \quad (\text{absorption})$$

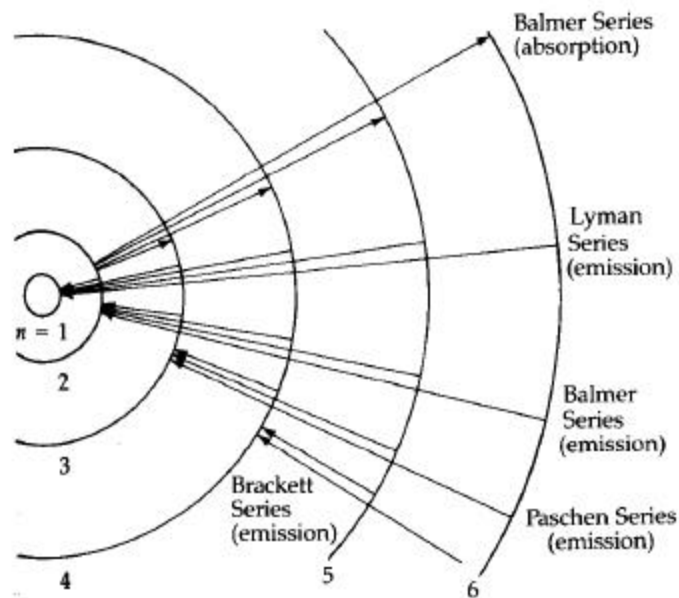
Bohr Model of the Hydrogen Atom

$$Z = 1 \quad \Rightarrow \quad E(n) = - (2p^2 m e^4 k^2) / n^2 h^2.$$

For a transition between level n_a and n_b :

$$1/\lambda_{ab} = R(1/n_b^2 - 1/n_a^2), \quad \text{where } \lambda_{ab} = \text{photon wavelength and}$$

$$R = \text{Rydberg constant} = 10.96776 \mu\text{m}^{-1}.$$



Note Lyman, Balmer, Paschen, Brackett, Pfund series for $n_b=1,2,3,4,5$.

Only Balmer series falls in visible light range, e.g., H α (Balmer α) line

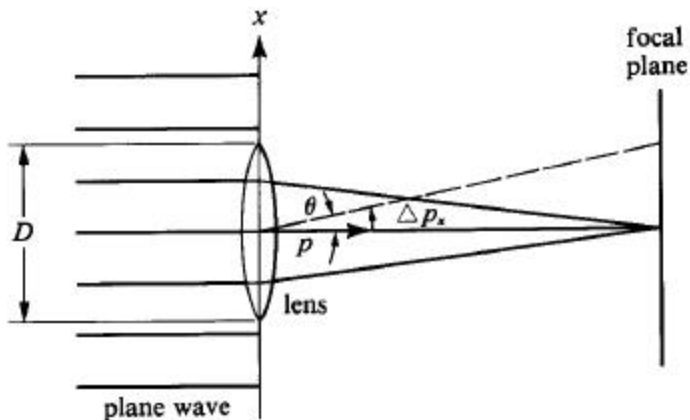
$$\lambda_{32} = 653.6 \text{ nm}.$$

Back to Wave-Particle Duality

No well-defined orbit radii. Logical conclusion of wave-particle duality \Rightarrow electrons do not have any definite position.

Heisenberg uncertainty principle: $\Delta x \Delta p > h$.

Applies to particles and waves. Explains diffraction and interference of particles.

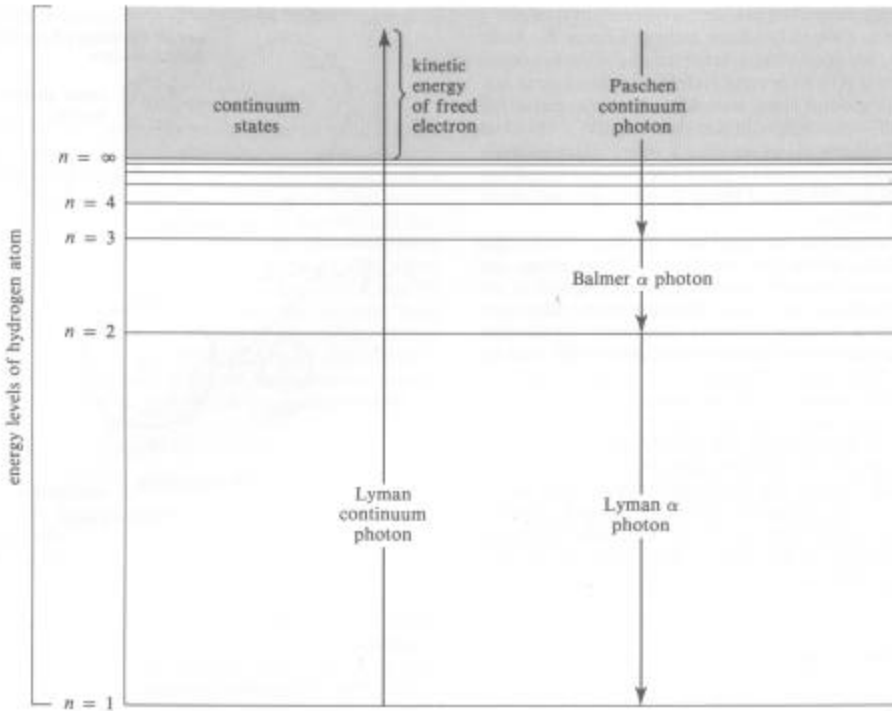


$$\Delta p > \frac{h}{D} \quad \text{and} \quad p = \frac{h}{\lambda} \quad \Rightarrow \quad \frac{\Delta p}{p} > \frac{\lambda}{D}.$$

$$\text{For } \Delta p \ll p, \quad \mathbf{q} \cong \frac{\Delta p}{p} = \frac{\lambda}{D}.$$

A Modern View of the Hydrogen Atom

Modern calculations of atomic structure still have discrete energy levels, if not well-defined electron radii. Electron radii described by a “wave” of probability around proton. Bohr orbits are most probable electron positions.



$$E_n = -\frac{13.6 \text{ eV}}{n^2}.$$

A Modern View of the Hydrogen Atom

In general, the energy of an electron depends on a host of quantum numbers:

Principal quantum number $n = 1, 2, 3, \dots$

Angular momentum quantum number $l = 0, 1, 2, \dots, n-1$

Angular momentum orientation $m_l = -l, -l+1, \dots, 0, \dots, l-1, l$

Spin $s = 1/2$

Spin orientation $m_s = -1/2$ or $1/2$.

Pauli Exclusion Principle: No two electrons may have an identical set of quantum numbers \Rightarrow all electrons cannot occupy the lowest energy state.

Multi-Electron Atoms

Transitions more complex than for hydrogen atom.

Electrons grouped into “shells”. If outermost shell is not full, electrons there easily excited and/or ionized. If only one electron in outermost shell, transitions similar to H.

Ions: when one or more electrons escape to the continuum. If one electron remaining, energy levels similar to H.

$$E_n = -\frac{13.6 \text{ eV } Z^2}{n^2}, \quad \text{e.g., He II, Li III, O VIII, Fe XXVI.}$$

Nomenclature: He⁺=He II, Ca = Ca I, Ca⁺=Ca II, Ca⁺⁺=Ca III

Molecules

Complex states.

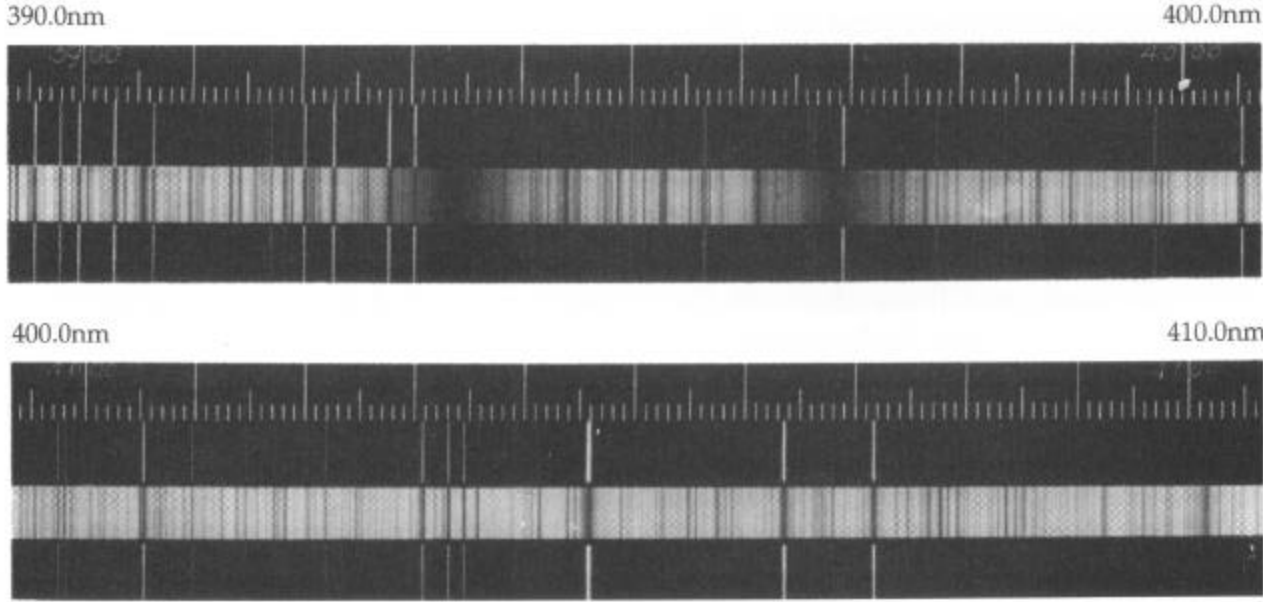
Transitions: electronic, rotational, vibrational.

For example, CO rotational levels

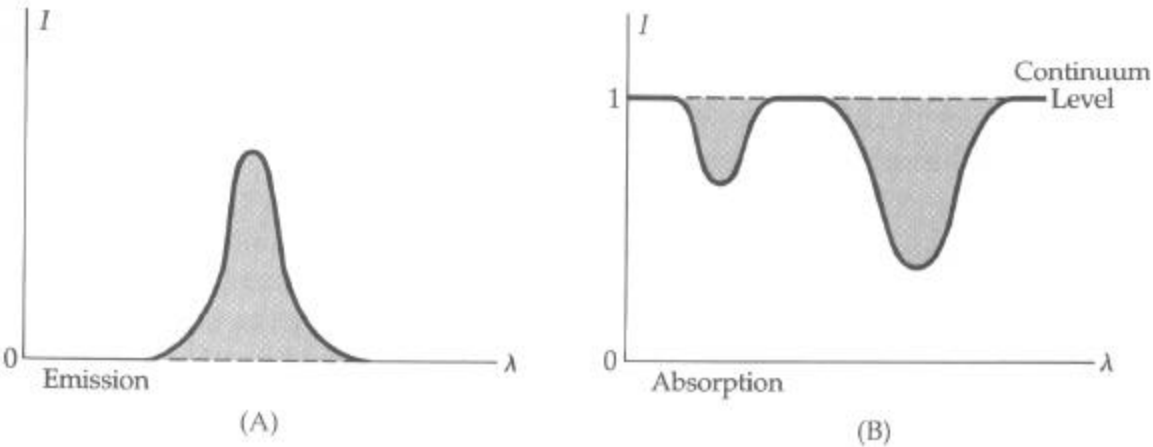
$$E = \left(\frac{h}{2\mathbf{p}} \right)^2 \frac{J(J+1)}{2\mathbf{m}r^2}, \quad J = 0, 1, 2, \dots$$

$J: 1 \rightarrow 0$ $\mathbf{n} = 115.3 \text{ GHz}$, $\mathbf{l} = 2.6 \text{ mm}$.

Spectral Lines



Solar spectrum



Emission and absorption lines

Spectral Lines

Absorption line: excitation to higher energy level by absorbing photon from background continuum with energy $\Delta E = h\nu$.

Subsequently, have

- (1) de-excitation by collision, so no photon produced
- (2) de-excitation to a different level, so photon of different energy (and ν) produced
- (3) de-excitation to original level, but photon (of same energy as original photon) may be emitted in any direction.

Emission line: Collisional excitation to higher level, followed by radiative decay so that a photon of energy $\Delta E = h\nu$ produced.

Forbidden emission line: Transition would be collisionally de-excited in laboratory, but de-excitation occurs radiatively in rarified interstellar gas.

Spectral Line Broadening

(1) natural broadening

$$\Delta E \Delta t > \frac{h}{2p} \Rightarrow \text{a natural width to energy levels}$$

$$\text{Lifetime of excited level } \Delta t \sim 10^{-8} \text{ s} \Rightarrow \Delta n = \frac{\Delta E}{h} = \frac{1}{2p\Delta t}.$$

$$c = n\mathbf{l} \Rightarrow \Delta \mathbf{l} = \frac{\mathbf{l}^2}{c} \Delta n = 1.3 \times 10^{-5} \text{ nm for } \mathbf{l} = 500 \text{ nm}.$$

(2) collisional broadening

– energy levels shifted by electrostatic interaction with neighboring atoms

– a direct dependence on particle density

– in a gas of sufficient density, characteristic spectral features disappear; gas emits a continuum of wavelengths.

Spectral Line Broadening

(3) Zeeman effect

A splitting of energy levels due to magnetic (**B**) field.

If splitting resolved, measure **B** strength.

If splitting unresolved, see broadened line.

(4) thermal Doppler broadening

Atomic motions along line-of-sight => Doppler shifts

$$\frac{\Delta l}{l} = \frac{v}{c}, \quad \text{e.g., for H at } T = 6000 \text{ K,}$$

$$\langle v \rangle \approx 12 \text{ km/s} \quad \Rightarrow \quad \Delta l \approx 0.025 \text{ nm}$$

for $l = 653.6 \text{ nm}$ **H α** line.

Spectral Line Broadening

(5) macroscopic broadening

Due to Doppler shifts from large-scale unresolved motions

- turbulence
- expansion or contraction
- rotation

Kirchoff's Rules

Relate appearance of spectra to the composition and physical state of an object.

(1) A hot and opaque solid, liquid, or highly compressed gas emits a continuous spectrum.

(2) A hot, transparent gas produces a spectrum of emission lines. Specific lines depend on which elements present in gas.

(3) Relatively cool, transparent gas in front of a continuum source produces absorption lines. Specific lines depend on which elements present in gas.